

RAPID ADIABATIC PREPARATION OF INJECTIVE PEPS AND GIBBS STATES

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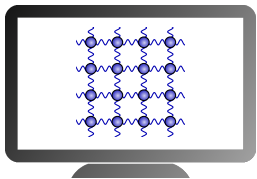
Yimin Ge, András Molnár, Ignacio Cirac

QIP 2016, 15.01.2016

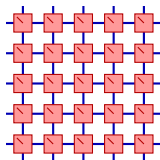
Max Planck Institute of Quantum Optics, Garching, Germany

MOTIVATION

Quantum simulation

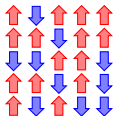


Quantum many-body systems



State preparation algorithm

Classical Gibbs distributions

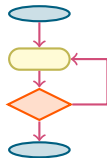


Statistical physics

Machine learning

Optimisation problems

Quantum algorithms



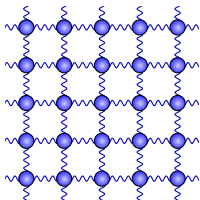


1. Introduction
 - PEPS
 - Gibbs states
 - Parent Hamiltonians
2. Naive algorithm
 - Adiabatic theorem
 - Hamiltonian simulation
3. Algorithm
 - Improved adiabatic theorem
 - Local changes
 - Lieb-Robinson localisation
4. Runtime & gap assumption

INTRODUCTION

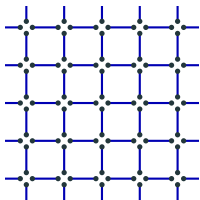
PROJECTED ENTANGLED PAIR STATES - PEPS

- Quantum system with N particles - $\dim \mathcal{H} = d^N$
 - Too large to describe efficiently
- Ground states of local gapped Hamiltonians
 - Special entanglement properties \rightarrow better description?



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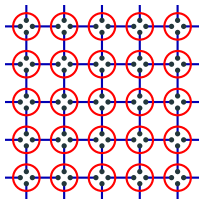
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$$\bullet \bullet = \sum_{i=1}^D |ii\rangle$$

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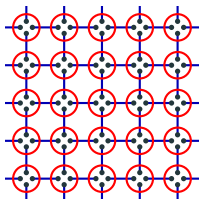
$$\text{---} \bullet = \sum_{i=1}^D |ii\rangle$$

$$\bigcirc = Q_\lambda : (\mathbb{C}^D)^{\otimes 4} \rightarrow \mathbb{C}^d$$

$$|\psi_{\text{PEPS}}\rangle = \bigotimes_{\text{sites } \lambda} Q_\lambda \bigotimes_{\text{edges}} \text{---} \bullet$$

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Injective PEPS

Q_λ invertible $\forall \lambda$

(generic case)

Gibbs state

- $H = \sum_{\mu} h_{\mu} \quad [h_{\mu}, h_{\nu}] = 0$
- $\rho_{\beta} = \frac{1}{Z} e^{-\beta H} \quad Z = \text{Tr} e^{-\beta H}$
- $\beta = \text{Temperature}^{-1}$

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Purification

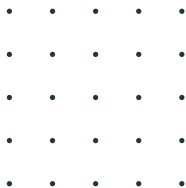
- $|\Psi\rangle = \frac{1}{\sqrt{Z}} (e^{-\beta H/2} \otimes \mathbb{1}) |\phi\rangle^{\otimes N}$
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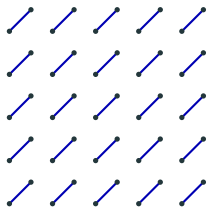
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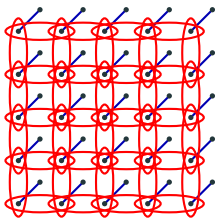
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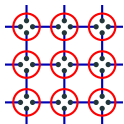
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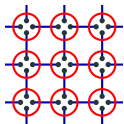
$$\text{---}\circ = e^{-\beta h_{\mu}/2}$$

Injective PEPS

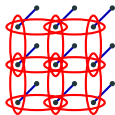


GENERAL CLASS OF STATES - PARENT HAMILTONIANS

Injective PEPS

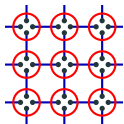


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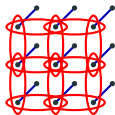


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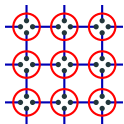


General

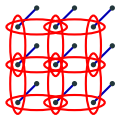


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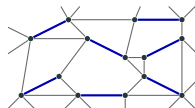
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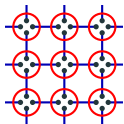


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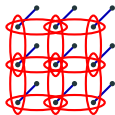


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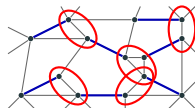
Injective PEPS



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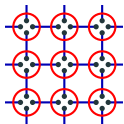


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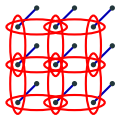


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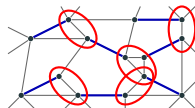
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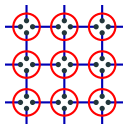
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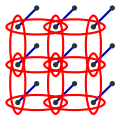
Commuting, finite range, invertible operators Q_λ acting on maximally entangled pairs

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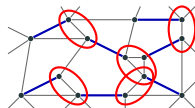
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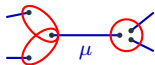


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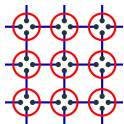
Parent Hamiltonians



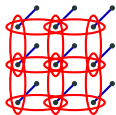
$|\Psi\rangle$

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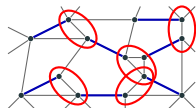
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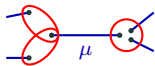


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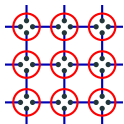
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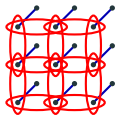
$$\left(\prod_{\lambda: \lambda \cap \mu \neq \emptyset} Q_\lambda^{-1} \right) |\Psi\rangle$$

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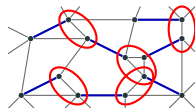
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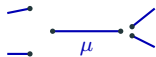


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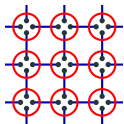
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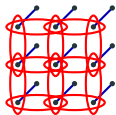
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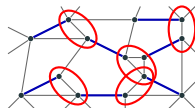
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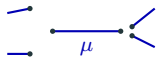


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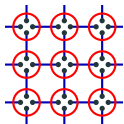
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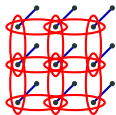
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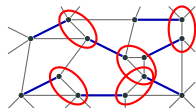
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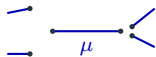


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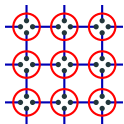
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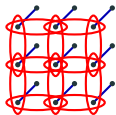
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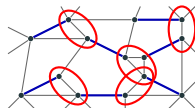
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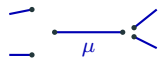


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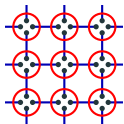
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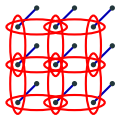
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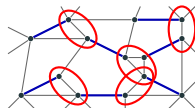
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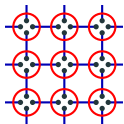
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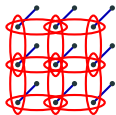
$$\sum_{\mu} \left(\prod_{\lambda: \lambda \cap \mu \neq \emptyset} Q_{\lambda}^{-1} \right)^{\dagger} P_{\mu} \left(\prod_{\lambda: \lambda \cap \mu \neq \emptyset} Q_{\lambda}^{-1} \right) |\Psi\rangle = 0$$

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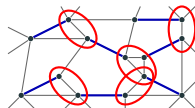
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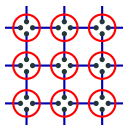
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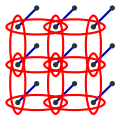
$$G = \sum_{\mu} \left(\prod_{\lambda: \lambda \cap \mu \neq \emptyset} Q_{\lambda}^{-1} \right)^{\dagger} P_{\mu} \left(\prod_{\lambda: \lambda \cap \mu \neq \emptyset} Q_{\lambda}^{-1} \right)$$

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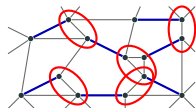
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Parent Hamiltonians

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$|\Psi\rangle$ is **unique** ground state of G

A NAIVE ALGORITHM

Parent Hamiltonian

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Adiabatic path

$$G(s) = \sum_{\mu} G_{\mu}(s), \quad s \in [0, 1]$$

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Adiabatic path

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Adiabatic theorem

e.g. [Jansen, Ruskai, Seiler '07]

Hamiltonian path $G(s)$ - ground state $|\phi(s)\rangle$ $s \in [0, 1]$ Initial state $|\psi(0)\rangle = |\phi(0)\rangle$ $\Delta = \min_s \text{Gap } G(s)$

$$i \frac{d}{dt} |\psi(t)\rangle = G\left(\frac{t}{\tau}\right) |\psi(t)\rangle \quad t \in [0, \tau]$$

$\tau = O\left(\frac{\|\dot{G}\|^2}{\varepsilon \Delta^3}\right)$ sufficient for final error $\|\psi(\tau) - \phi(1)\| < \varepsilon$

Parent Hamiltonian

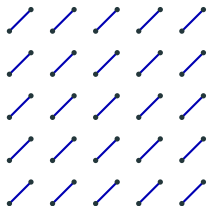
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Adiabatic path

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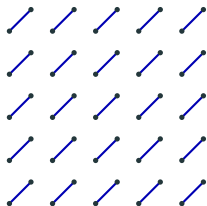


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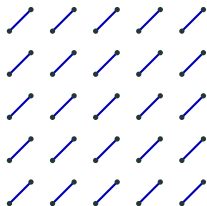
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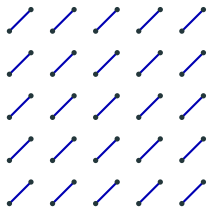
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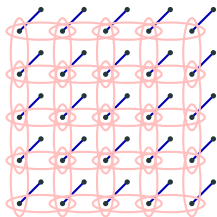
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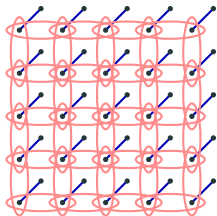
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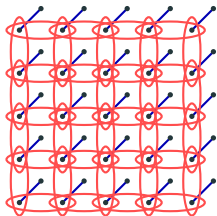
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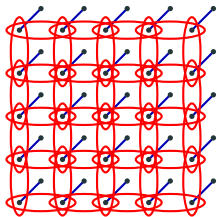
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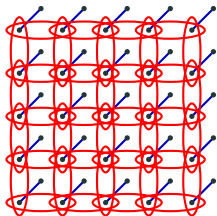
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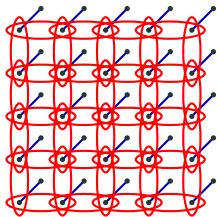
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Adiabatic runtime $\tau = O\left(\frac{\|\dot{G}\|^2}{\varepsilon \Delta^3}\right)$

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Hamiltonian simulation

e.g. [Berry, et al '14]

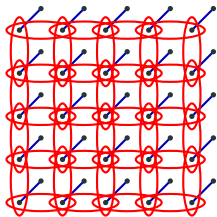
Given: Hamiltonian $G(t)$ $t \in [0, \tau]$

Want: Unitary evolution of $G \rightarrow$ quantum circuit

$$\Rightarrow \# \text{ gates} = \tilde{O}((\# \text{ qubits}) \times \|G\| \times \tau)$$

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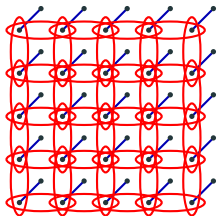
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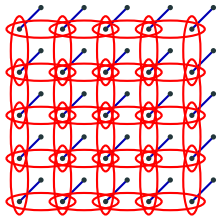
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Runtime (# gates) $T = \tilde{O}(\tau N^2)$

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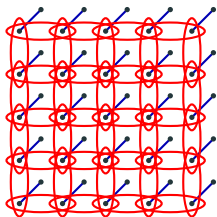
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Runtime (# gates) $T = \tilde{O}(N^4 \Delta^{-3} \epsilon^{-1})$

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Worse than classical algorithms
for classical Gibbs states



IMPROVED ALGORITHM

Summary of results

Algorithm with

Runtime (# gates) $T = O(N \text{ polylog}(N/\epsilon))$ ← almost optimal

Circuit depth $D = O(\text{polylog}(N/\epsilon))$

assuming a *uniform gap*

Ingredients

1. Adiabatic theorem with (almost) exponentially small error
2. Sequence of local changes
3. Lieb-Robinson localisation

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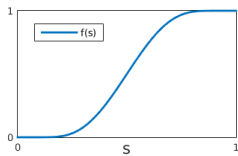
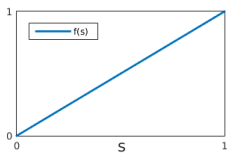


Main message: **Speedup by locality**

ADIABATIC THEOREM WITH ALMOST EXPONENTIALLY SMALL ERROR

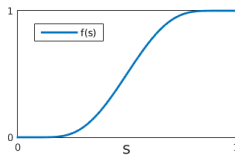
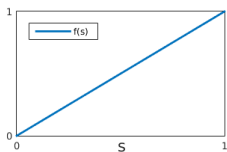
Smooth reparameterisation

$$G(s) \rightarrow G(f(s))$$



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Theorem

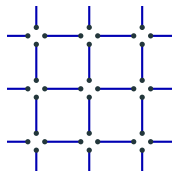
c.f. [Nenciu '93], [Hagedorn, Joye '02]

- $G(s)$ smooth & all derivatives of $G(s)$ vanish at $s = 0, 1$
- $G(s)$ is in Gevrey class $1 + \alpha$
- $K = |\text{supp } \dot{G}|$

$$\Rightarrow \tau = O\left(\log^{1+\alpha}\left(\frac{K}{\varepsilon\Delta}\right) \frac{K^2}{\Delta^3}\right) \quad \text{sufficient for final error } \varepsilon$$

SEQUENCE OF LOCAL CHANGES

$$G_{\text{initial}} = \sum_{\mu} P_{\mu}$$

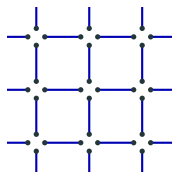


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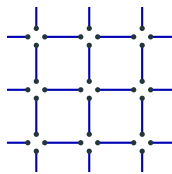
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Change Q_{λ} 's
individually

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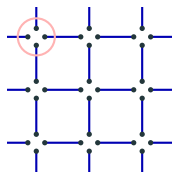
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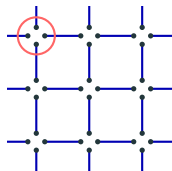
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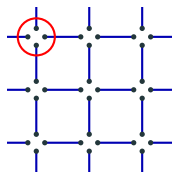
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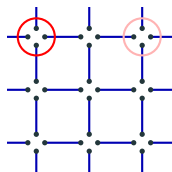
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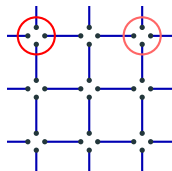
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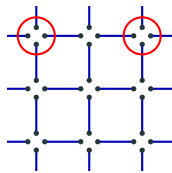
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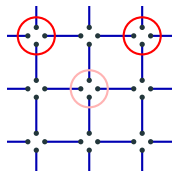
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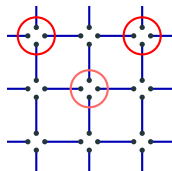
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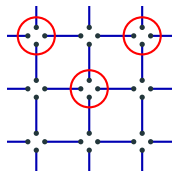
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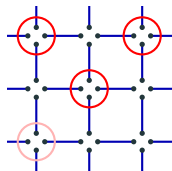
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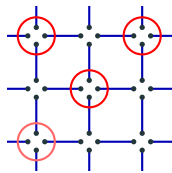
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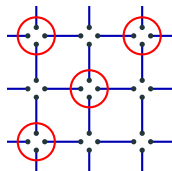
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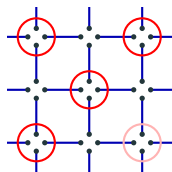
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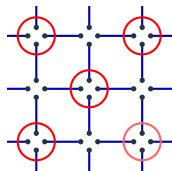
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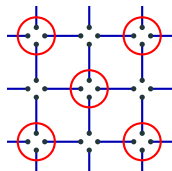
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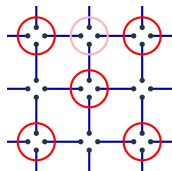
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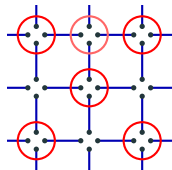
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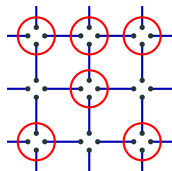
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Change Q_{λ} 's
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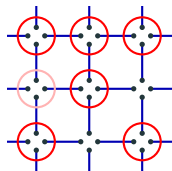
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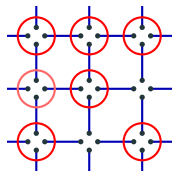
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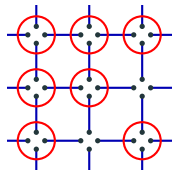
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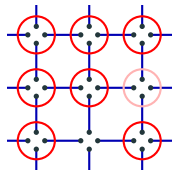
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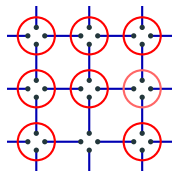
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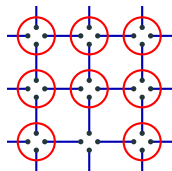
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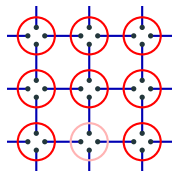
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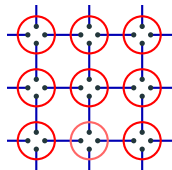
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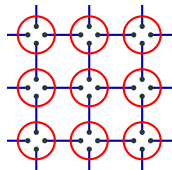
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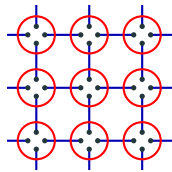
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Δ = minimum gap

Runtime

Adiabatic runtime N paths $\times O\left(\log^{1+\alpha}(N/\epsilon\Delta)\Delta^{-3}\right)$

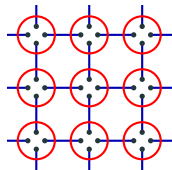
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Adiabatic runtime $N \text{ paths} \times O\left(\log^{1+\alpha}(N/\varepsilon\Delta)\Delta^{-3}\right)$

Runtime (# gates) $T = O\left(N^3 \text{ polylog}(N/\varepsilon\Delta)\Delta^{-3}\right)$

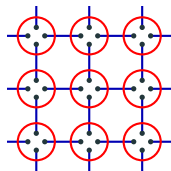
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Change Q_{λ} 's individually

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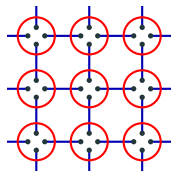
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Problem: Hamiltonians act on **entire** system!

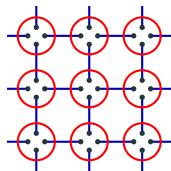
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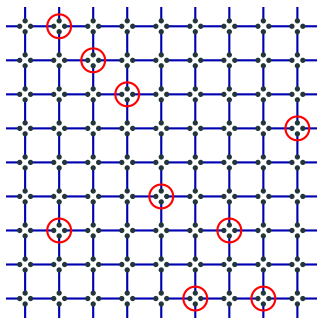
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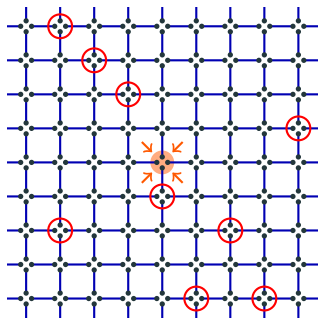
Problem: Hamiltonians act on **entire** system!

Assume from now: $\Delta = \Omega(1)$

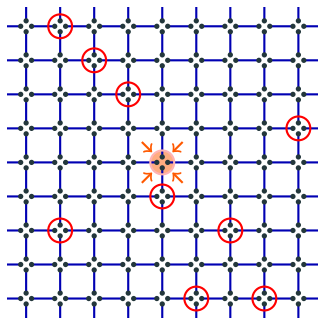
LIEB-ROBINSON LOCALISATION



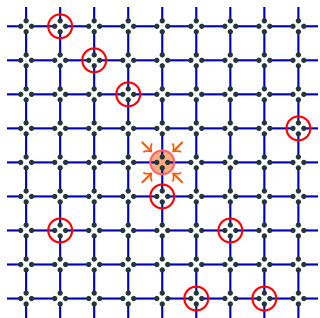
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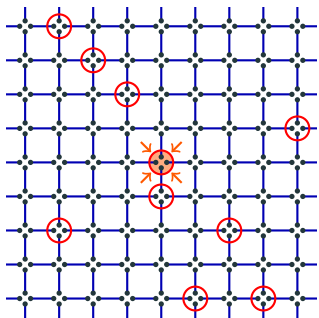
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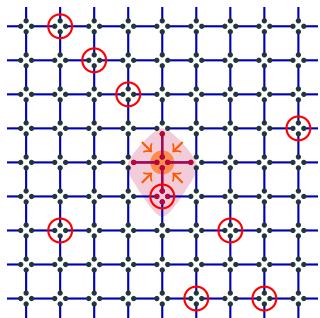
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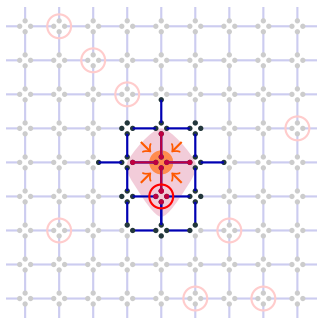


LIEB-ROBINSON LOCALISATION



supp \dot{G}

LIEB-ROBINSON LOCALISATION



Idea

Localise Hamiltonian around $\text{supp } \dot{G}$

Tool

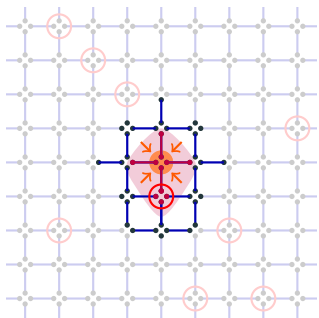
Lieb-Robinson bound

Assumptions

$$\tau = O(\text{polylog}(N/\varepsilon))$$

G frustration-free

LIEB-ROBINSON LOCALISATION



Idea

Localise Hamiltonian around $\text{supp } \dot{G}$

Tool

Lieb-Robinson bound

Assumptions

$\tau = O(\text{polylog}(N/\epsilon))$

G frustration-free

Theorem

Hamiltonian terms supported at distance $\gtrsim \text{polylog}(\frac{N}{\epsilon})$ away from $\text{supp } \dot{G}$ have negligible effect on adiabatic evolution

Truncate \rightarrow Hamiltonians act only on $O(\text{polylog}(\frac{N}{\epsilon}))$ sites

RUNTIME & GAP ASSUMPTION

Runtime

Adiabatic runtime:

$$N \text{ paths} \times O\left(\log^{1+\alpha}\left(\frac{N}{\epsilon}\right)\right)$$

Each path:

$$\text{supported on } O\left(\text{polylog}\left(\frac{N}{\epsilon}\right)\right) \text{ sites}$$



Runtime

- Adiabatic runtime: $N \text{ paths} \times O\left(\log^{1+\alpha}\left(\frac{N}{\epsilon}\right)\right)$
- Each path: supported on $O\left(\text{polylog}\left(\frac{N}{\epsilon}\right)\right)$ sites
- Runtime (# gates): $T = O\left(N \text{ polylog}\left(\frac{N}{\epsilon}\right)\right)$



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Circuit depth: $D = O\left(\text{polylog}\left(\frac{N}{\epsilon}\right)\right)$

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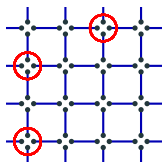
Assumption

$\Omega(1)$ gap along all N paths



UNIFORM GAP ASSUMPTION

G_{n-1}
↓
change
single Q_λ
↓
 G_n



$$\bigcirc = Q_\lambda : (\mathbb{C}^D)^{\otimes 4} \rightarrow \mathbb{C}^d$$

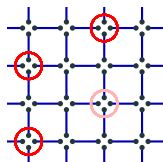
Assumption

$$\text{Gap } G_n(s) \geq \text{const} > 0$$

$$\forall n = 1 \dots N \text{ and } s \in [0, 1]$$

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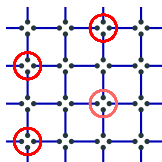
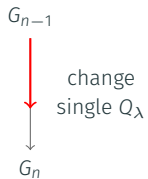
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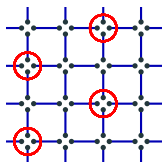
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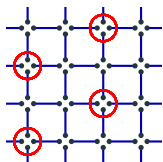
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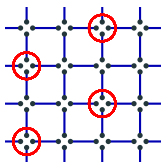
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If $Q_\lambda > q\mathbb{1}$ for all λ and some $q = \text{const} > 0$

$$\Rightarrow \text{Gap } G_n(s) \geq q^{-2} \text{Gap } G_n(0) \quad \forall s \in [0, 1].$$

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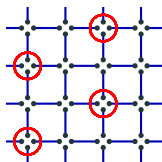
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Uniform gap assumption

$$\text{Gap } G_n(0) = \Omega(1) \quad \forall n = 1 \dots N$$

“Parent Hamiltonian of various system sizes (& shapes) gapped”

COMPARISON WITH OTHER ALGORITHMS



Runtime
(uniform gap)

Our algorithm

$O(N \text{ polylog } N)$

COMPARISON WITH OTHER ALGORITHMS



Runtime
(uniform gap)

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Injective PEPS

[Schwarz et al '12]

$O(N^4)$

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MCMC (e.g. Metropolis)

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Additional assumptions

$O(N \log N)$

← “rapid mixing”

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Quantum

QSA [Somma et al '08]

[Wocjan, Abeyesinghe '08]

QM [Temme et al '11]

$O(\text{poly } N)$

QQM [Yung, Aspuru-Guzik '12]

...

COMPARISON WITH OTHER ALGORITHMS



Runtime
(uniform gap)

Gap
dependence

Our algorithm

$O(N \text{ polylog } N)$

$\tilde{O}(\Delta^{-3})$ or $\tilde{O}(\Delta^{-3-6 \dim})$

Injective PEPS

[Schwarz et al '12]

$O(N^4)$

$O(\Delta^{-1})$

Gibbs states

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MCMC (e.g. Metropolis)

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[Wocjan, Abeyesinghe '08]

$O(\Delta^{-1/2})$

QM [Temme et al '11]

$O(\text{poly } N)$

$O(\Delta^{-1})$

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- Sequence of local changes
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