On the Power of Quantum Memory

Ueli Maurer, ETH Zurich

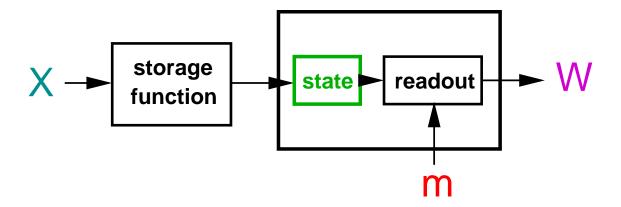
Joint work with Robert König and Renato Renner

paper at: quant-ph/0305154

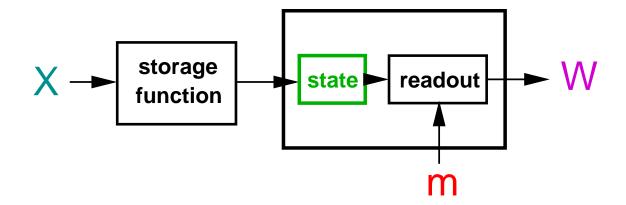
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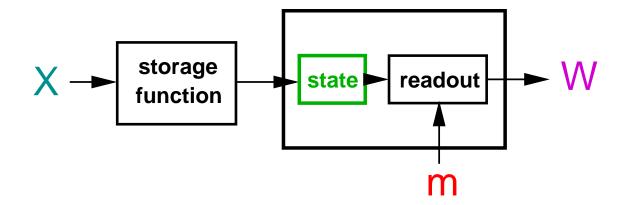


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- Is privacy amplification secure against an adversary holding quantum information?
- Christandel's talk: Implications to quantum cryptography?

Overview

1. Information-theoretic cryptography

2. Characterizing the power of quantum storage

3. Privacy amplification is secure against quantum adversaries

- Randomness exists (generation of secret keys)
- Independence exists (
 ∃ telepathy)

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- Computational intractability assumptions

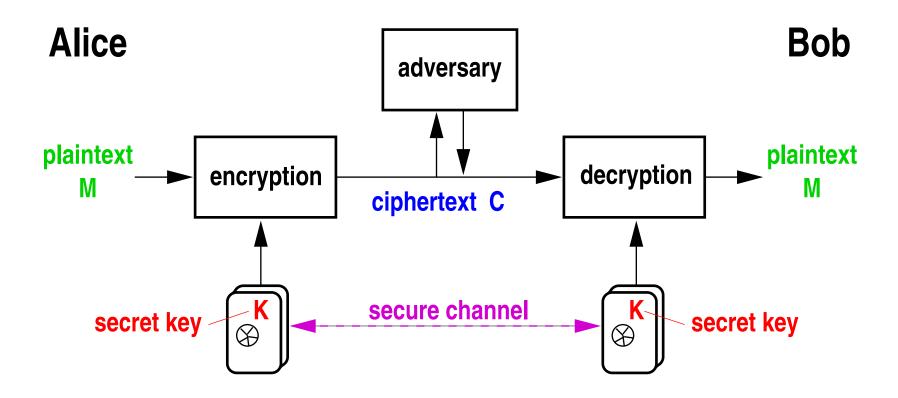
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- Physical assumptions
 - Tamper-resistance
 - Noise in communication systems
 - Restrictions on adversary's memory capacity
 - Quantum theory

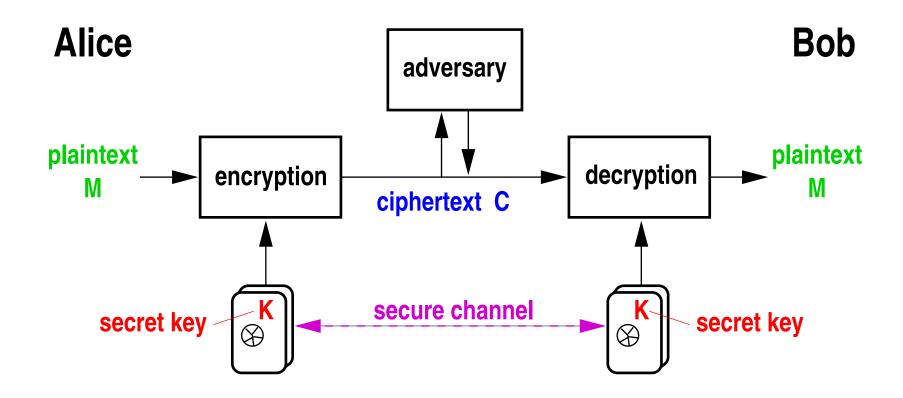
Why cryptography without comp. assumptions

- Which is the right model of computation?
- No lower bound proofs for any useful comput. model.
- Clean security definitions.
- Physical assumptions are more sound than comp. ass.

Symmetric cryptosystem

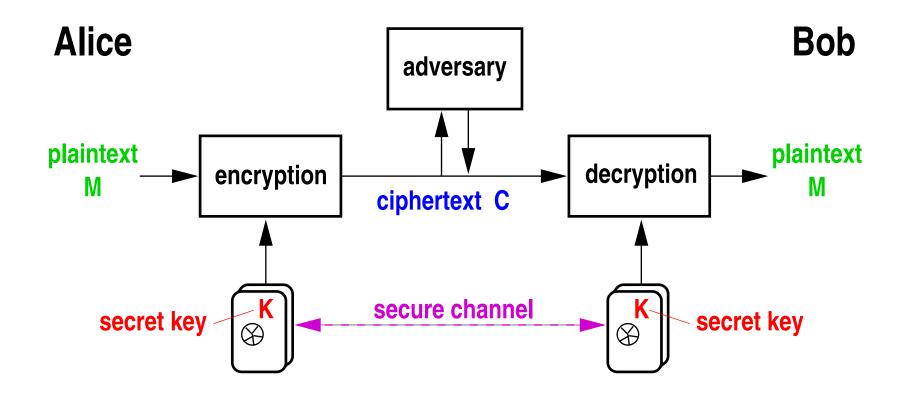


Symmetric cryptosystem



Definition: A cryptosystem is perfect if I(M;C) = 0.

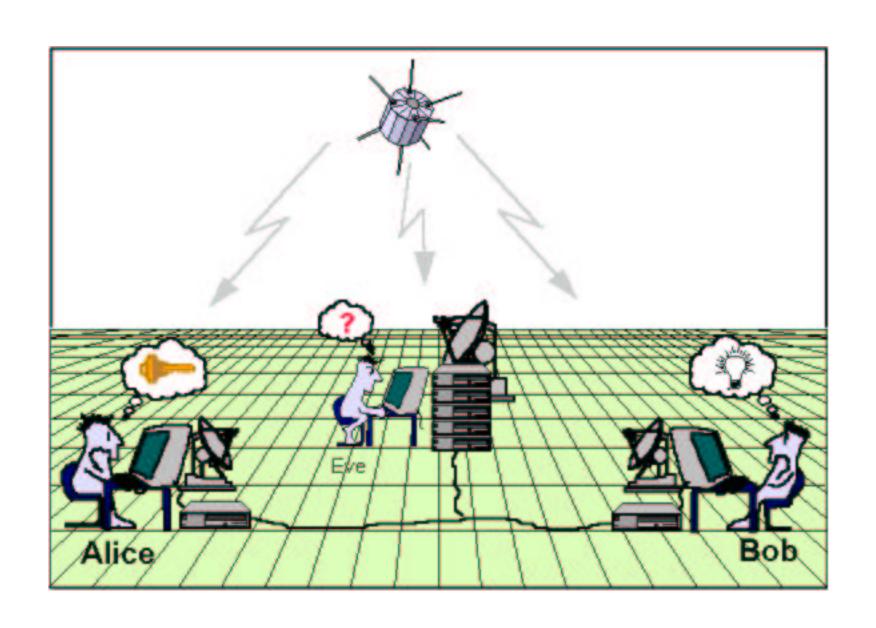
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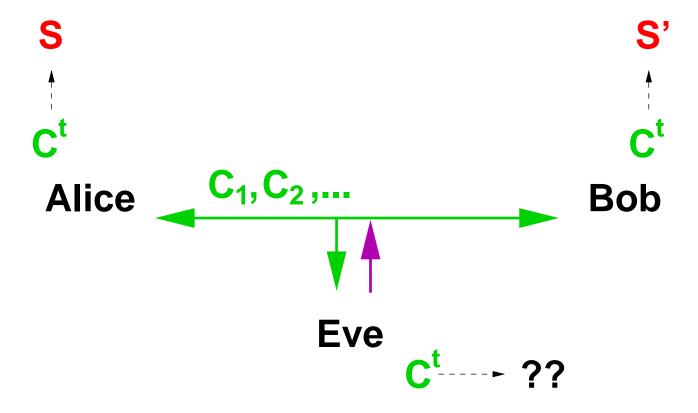


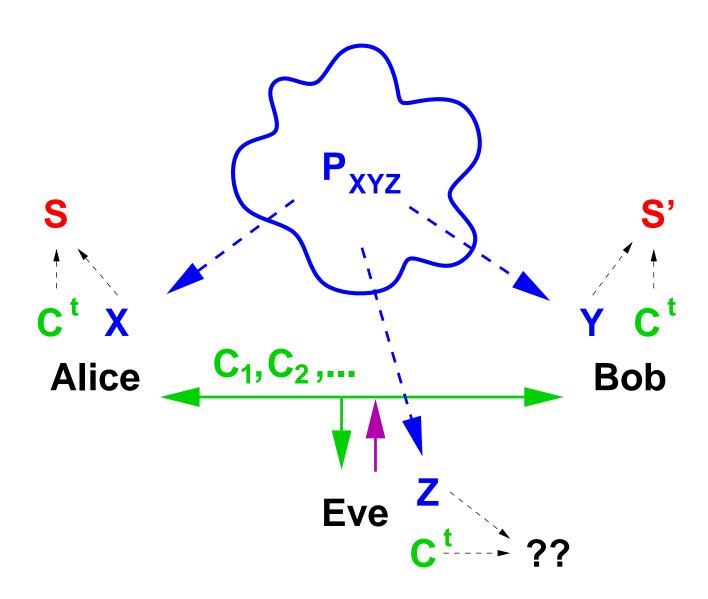
Definition: A cryptosystem is perfect if I(M;C) = 0.

Theorem [Sha49]: For every perfect cipher, $H(K) \ge H(M)$.

Information-theoretic key agreement by public discussion [M93]







Corollary: $H(K) \ge H(M)$ holds also in an interactive setting.

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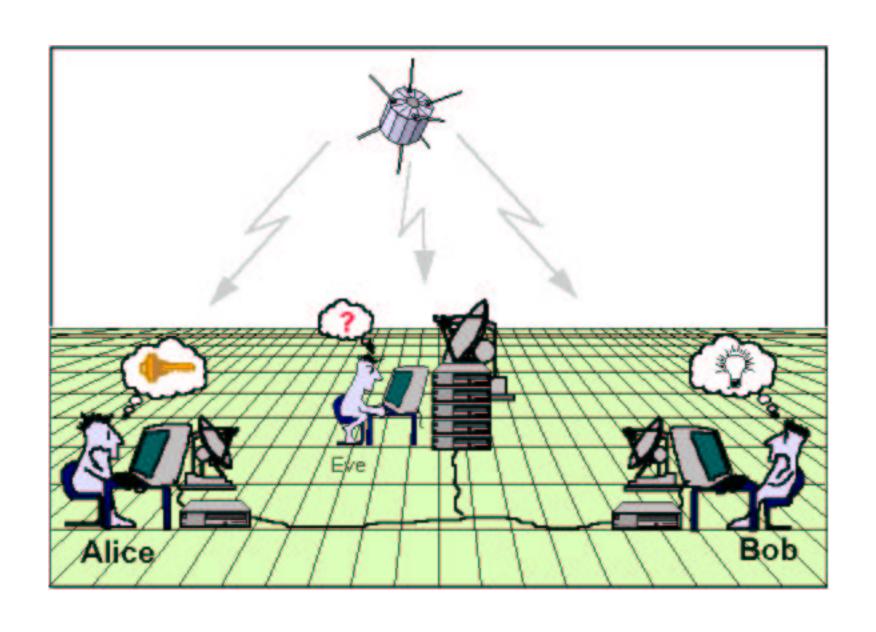
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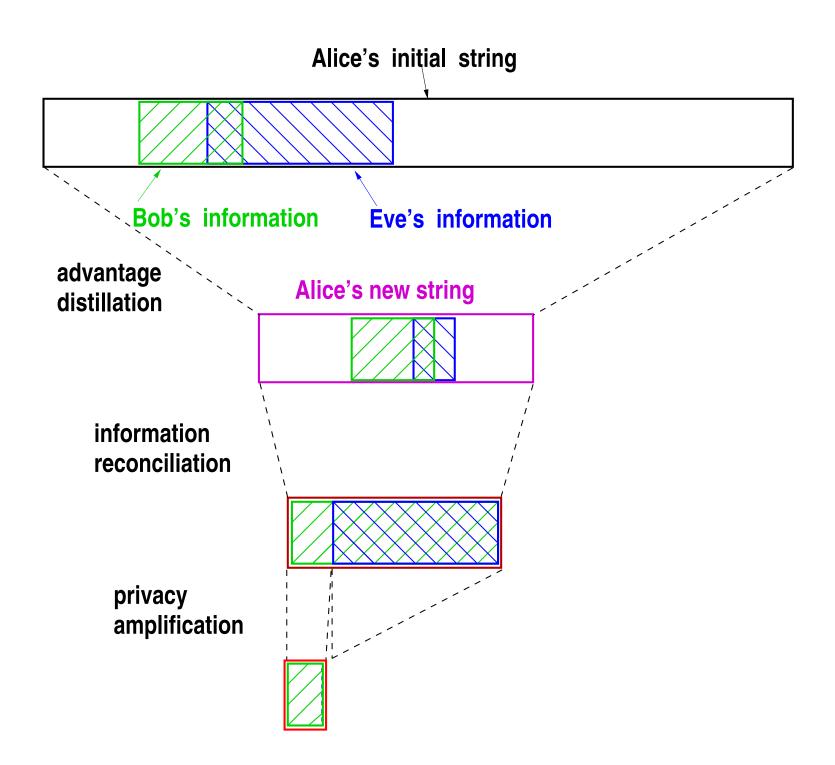
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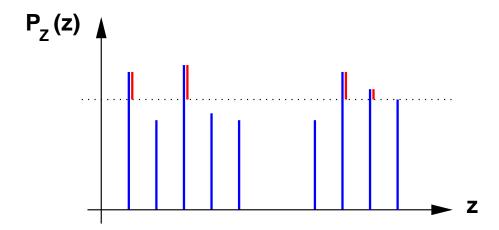
Theorem: In the satellite model, H(S) > 0 is possible whenever it is not obviously impossible, i.e., if

- Eve's channel is not perfectly noiseless and
- Alice's and Bob's channels have positive capacity.

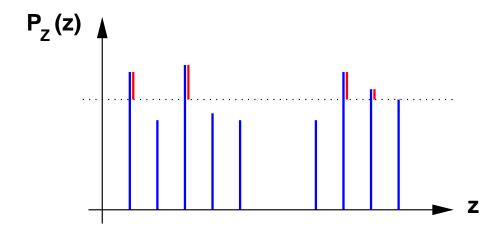
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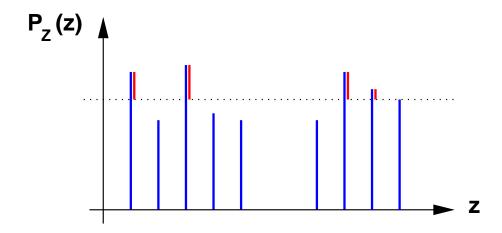


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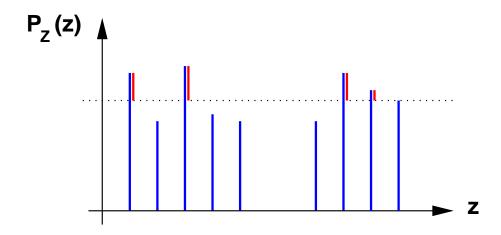
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Lemma: One can define a uniform random variable **Z** that is independent of W and such that **Z** = **Z** holds with probability $1 - d(\mathbf{Z}|\mathbf{W})$.



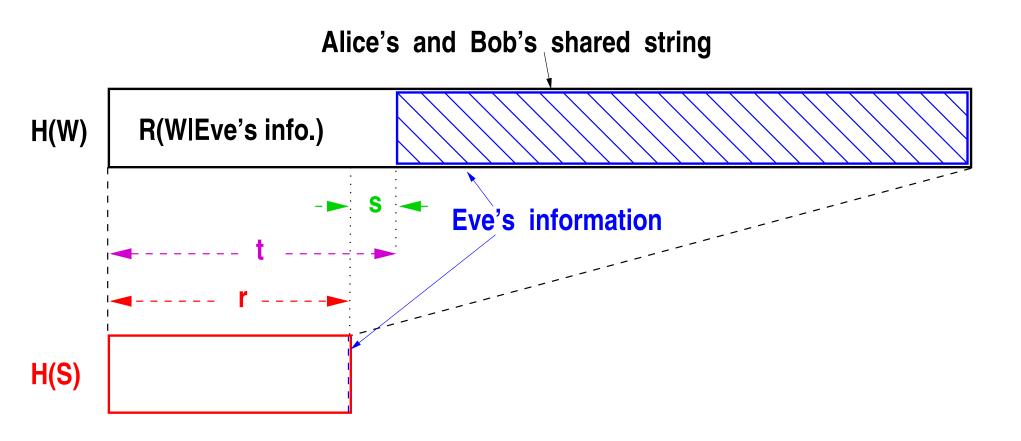
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In other words, with probability $1 - d(\mathbf{Z}|\mathbf{W})$ the setting with \mathbf{W} and \mathbf{Z} is equivalent to an ideal setting with \mathbf{W} and independent uniform \mathbf{Z} .

Privacy amplification



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Theorem: Let X and W be arbitrary random variable with $H_2(X|W) \ge t$ and let G be a 2-universal random function from \mathcal{X} to $\{0,1\}^s$. Then

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Corollary: If X is uniform over $\{0,1\}^n$ and W consists of r arbitrary (classical) bits about X, then

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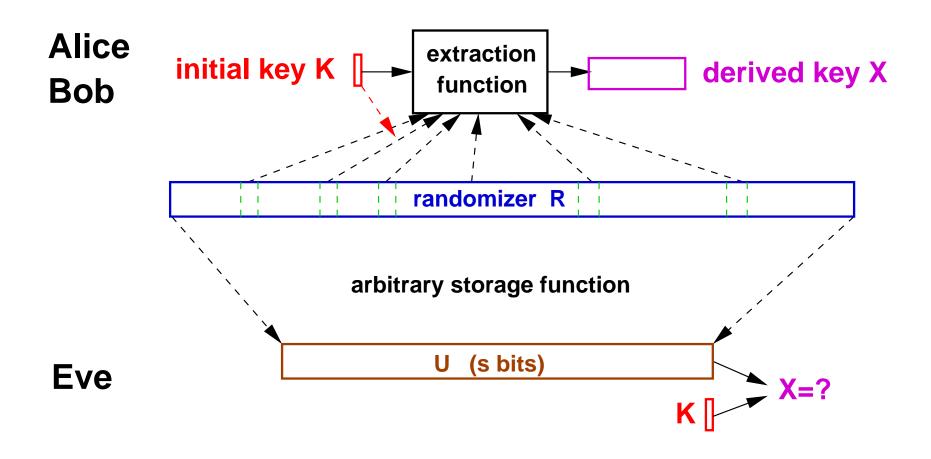
Question: What about quantum knowledge about X?

The bounded-storage model (BSM) [M90]

Basic idea: Eve has bounded storage capacity of s bits, but otherwise unlimited computing power.

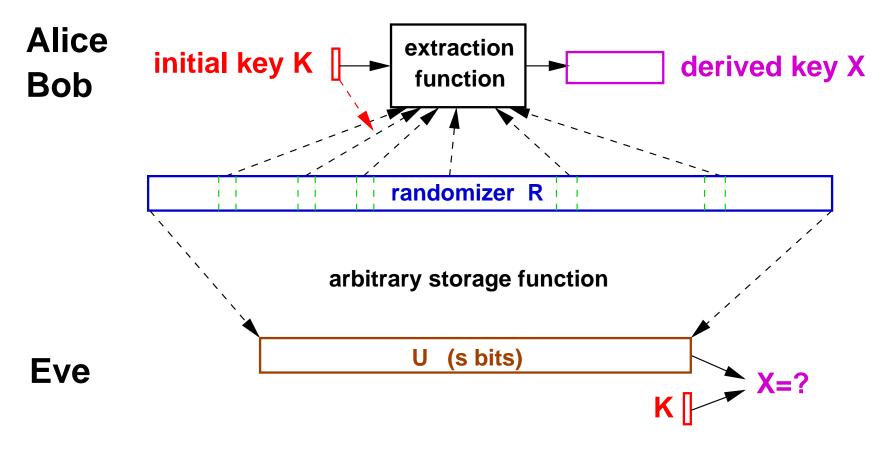
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Question: What about quantum storage?

Lemma: Consider any random variable Z over Z. If H is a uniform balanced Boolean random function, then

$$\mathbf{d}(\mathbf{Z}) \leq \frac{3}{2}\sqrt{|\mathcal{Z}|} \ \mathbf{d}(\mathbf{H}(\mathbf{Z})|\mathbf{H}).$$

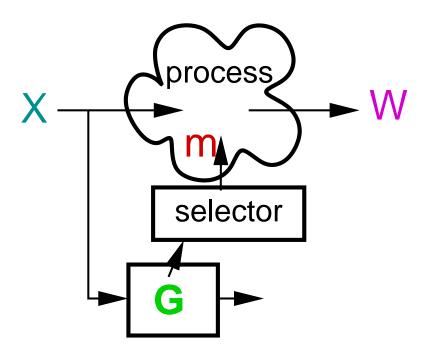
W

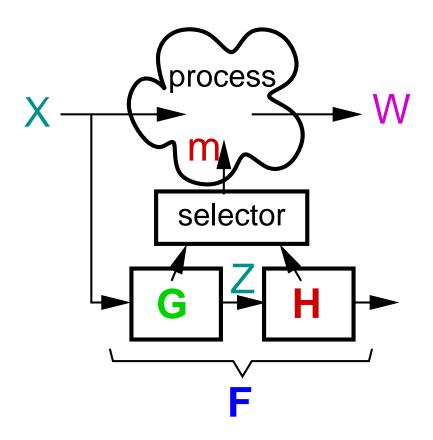
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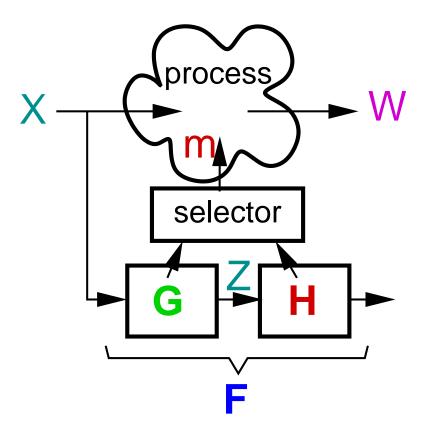
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More generally,

$$\mathbf{d}(\mathbf{Z}|\mathbf{W}) \leq \frac{3}{2}\sqrt{|\mathcal{Z}|} \ \mathbf{d}(\mathbf{H}(\mathbf{Z})|\mathbf{W}\mathbf{H}).$$

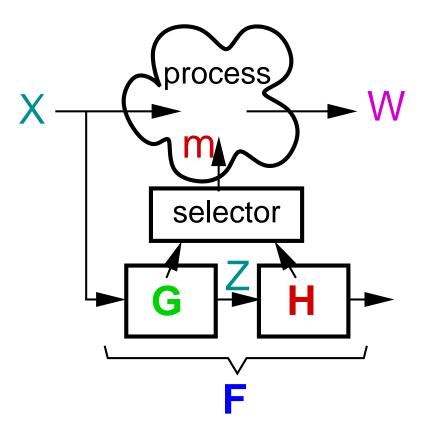






Corollary: Consider any process generating W from a random variable X and a selection input m. If for any 2-universal $(\mathcal{X}, \{0, 1\})$ -random function F and for any selector with input F we have

$$d(F(X)|WF) \leq \epsilon$$
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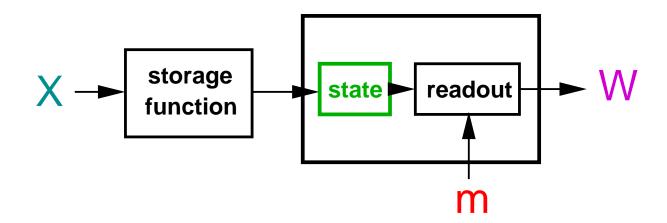


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then for any 2-universal $(\mathcal{X}, \{0, 1\}^s)$ -random function **G** and for any selector with input **G** $\mathsf{d}(\mathbf{G}(\mathbf{X})|\mathbf{WG}) \leq \tfrac{3}{2} \, 2^{s/2} \, \epsilon.$

r-qubit quantum storage device



State: Normalized vector ψ in the d-dimensional Hilbert space \mathcal{H}_d ($d=2^r$).

Equivalently, state space = $\mathcal{P}(\mathcal{H}_d) := \{P_\psi : \psi \in \mathcal{H}_d, \|\psi\| = 1\}$ (pure states), where P_ψ is the projection operator in \mathcal{H}_d along the vector ψ .

Most general read-out operation: $\mathbf{m} \in \mathsf{POVM}(\mathcal{H}_d)$, resulting in W.

m is specified by a family $\{E_w\}$ of nonneg. op. on \mathcal{H}_d with $\sum_w E_w = \mathrm{id}_{\mathcal{H}_d}$.

System in state $P_{\psi} \Rightarrow P_{\mathbf{W}}(w) = \operatorname{tr}(E_w P_{\psi})$.

The quantum binary decision problem

Given: A QS prepared in one of two mixed states $\rho_0, \rho_1 \in \mathcal{S}(\mathcal{H})$, with a priori probabilities q and 1-q, respectively.

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General decision strategy: POVM $\{E_0, E_1\}$

Prob[**W** =
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] = **tr** $(E_i \rho_j)$, for $i, j \in \{0, 1\}$.

Success probability: $q \operatorname{tr}(E_0 \rho_0) + (1-q) \operatorname{tr}(E_1 \rho_1)$.

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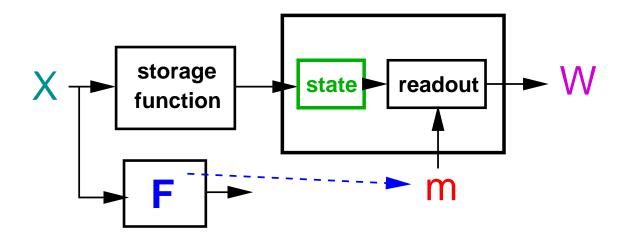
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Theorem [Hel76]: The maximum achievable success probability is

$$\frac{1}{2} + \frac{1}{2} \sum_{j=1}^{d} |\mu_j|$$
,

where $\{\mu_j\}_{j=1}^d$ are the eigenvalues of the hermitian operator

$$\Gamma := q \, \rho_0 - (1 - q) \, \rho_1.$$



Lemma: Let

 ${\sf X}$ = random variable with range ${\cal X}$, stored in an r-qubit quantum system using storage function $\varphi: x \mapsto P_{\psi_x}$.

 \mathbf{F} = any Boolean random function on \mathcal{X} .

W = measurement outcome of any measurement on the state,
 depending on F.

Then

$$\mathbf{d}(\mathbf{F}(\mathbf{X})|\mathbf{WF}) \leq \frac{1}{2} E_{\mathbf{F}} \left[\sum_{j=1}^{d} |\mu_{j}^{\mathbf{F}}| \right],$$

where for every f, $\{\mu_j^f\}_{j=1}^d$ are the eigenvalues of the hermitian operator

$$\Lambda_f := \sum_{x:f(x)=0} P_{\mathbf{X}}(x) P_{\psi_x} - \sum_{x:f(x)=1} P_{\mathbf{X}}(x) P_{\psi_x}$$

Let

$$\lambda_{x,x'} := 2 \operatorname{Prob}[F(x) = F(x')] - 1 = E_F[\delta_{f(x),f(x')} - 1]$$

Note: For 2-universal F, $\lambda_{x,x'} \leq 0$ for $x \neq x'$.

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Theorem: Let X, F, and W be as above. Then

$$\mathbf{d}(\mathbf{F}(\mathbf{X})|\mathbf{WF}) \ \leq \ \tfrac{1}{2} d^{\tfrac{1}{2}} \sqrt{\sum_{x,x' \in \mathcal{X}} P_{\mathbf{X}}(x) \, P_{\mathbf{X}}(x') \, \lambda_{x,x'}} \ \operatorname{tr}(P_{\psi_x} P_{\psi_{x'}})$$

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Corollary: If F is 2-universal, then

$$\mathbf{d}(\mathbf{F}(\mathbf{X})|\mathbf{WF}) \leq \frac{1}{2}d^{\frac{1}{2}}\sqrt{\sum_{x\in\mathcal{X}}P_{\mathbf{X}}^{2}(x)} = \frac{1}{2} 2^{\frac{1}{2}(H_{2}(\mathbf{X})-r)}$$

Moreover, if X is a uniform n-bit string, then

$$d(F(X)|WF) \leq \frac{1}{2} 2^{\frac{1}{2}(n-r)}$$

Theorem: Let X, F, and W be as above. Then

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Proof: For any f,

$$\sum_{j=1}^{d} |\mu_j^f| \leq d^{\frac{1}{2}} \sqrt{\sum_{j=1}^{d} |\mu_j^f|^2} = d^{\frac{1}{2}} \sqrt{\operatorname{tr}(\Lambda_f^2)},$$

(using Jensen's inequality and Schur's (in)equality).

$$\mathbf{d}(\mathbf{F}(\mathbf{X})|\mathbf{WF}) \ \leq \ \tfrac{1}{2} E_{\mathbf{F}} \Big[\sum_{j=1}^{d} |\mu_{j}^{\mathbf{F}}| \Big] \ \leq \ \tfrac{1}{2} \ d^{\frac{1}{2}} \ E_{\mathbf{F}} \Big[\sqrt{\mathbf{tr}(\Lambda_{\mathbf{F}}^{2})} \Big] \leq \tfrac{1}{2} \ d^{\frac{1}{2}} \sqrt{E_{\mathbf{F}}[\mathbf{tr}(\Lambda_{\mathbf{F}}^{2})]} \ .$$

$$\begin{aligned} \operatorname{tr}(\Lambda_f^2) &= \sum_{\substack{x,x' \in \mathcal{X} \\ f(x) = f(x')}} P_X(x) P_X(x') \operatorname{tr}(P_{\psi_x} P_{\psi_{x'}}) - \sum_{\substack{x,x' \in \mathcal{X} \\ f(x) \neq f(x')}} P_X(x) P_X(x') \operatorname{tr}(P_{\psi_x} P_{\psi_{x'}}) \\ &= \sum_{\substack{x,x' \in \mathcal{X} \\ x,x' \in \mathcal{X}}} 2(\delta_{f(x),f(x')} - 1) P_{\mathbf{X}}(x) P_{\mathbf{X}}(x') \operatorname{tr}(P_{\psi_x} P_{\psi_{x'}}) \\ &= \sum_{\substack{x,x' \in \mathcal{X} \\ x,x' \in \mathcal{X}}} 2(\delta_{f(x),f(x')} - 1) P_{\mathbf{X}}(x) P_{\mathbf{X}}(x') \operatorname{tr}(P_{\psi_x} P_{\psi_{x'}}) \end{aligned}$$

Comparing classical and quantum storage devices

Lemma: For a uniform 2-bit random variable X, a uniform Boolean balanced random function F, and a 1-(qu)bit storage system,

$$d_{\text{opt}}^{\mathsf{C}}(\mathbf{F}(\mathbf{X})|\mathbf{WF}) = \frac{1}{4}$$

and

$$d_{\mathrm{opt}}^{\mathsf{q}}(\mathbf{F}(\mathbf{X})|\mathbf{WF}) = \frac{1}{2\sqrt{3}} \approx 0.289$$
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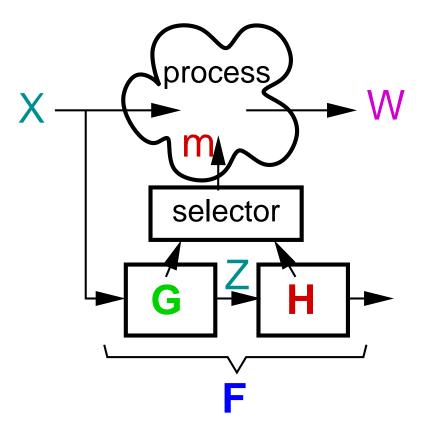
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Lemma: For any random variable X and any uniform random function F,

$$\frac{1}{\sqrt{2\pi}}(1+O(2^{-(n-r)}))2^{-\frac{n-r}{2}} \le d_{\text{opt}}^{\mathsf{C}}(\mathbf{F}(\mathbf{X})|\mathbf{WF}) \le d_{\text{opt}}^{\mathsf{q}}(\mathbf{F}(\mathbf{X})|\mathbf{WF}) \le \frac{1}{2}2^{-\frac{n-r}{2}}.$$



Corollary: Let X be a random variable over \mathcal{X} . If for any 2-universal $(\mathcal{X}, \{0, 1\})$ -random function F and for any process generating W from X and F we have

$$d(F(X)|WF) \leq \epsilon$$

then for any 2-universal $(\mathcal{X}, \{0, 1\}^s)$ -random function G and any process generating a random variable W from X and G we have

$$d(G(X)|WG) \leq \frac{3}{2} 2^{s/2} \epsilon$$
.

Privacy amplification is secure against quantum adversaries

Theorem: Let X be uniformly distributed over $\{0,1\}^n$ and let G be a 2-universal random function from $\{0,1\}^n$ to $\{0,1\}^s$. If all information about X is stored in r qubits, then

$$d_{\text{opt}}^{\mathsf{q}}(\mathbf{G}(\mathbf{X})|\mathbf{WG}) \leq \frac{3}{4} 2^{-\frac{1}{2}(n-r-s)}$$
.

Note:

$$d_{\text{opt}}^{\mathsf{C}}\left(\mathbf{G}(\mathbf{X})|\mathbf{WG}\right) = O\left(2^{-\frac{1}{2}(n-r-s)}\right)$$

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 In a quite general context quantum memory is only marginally more powerful than classical memory.

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- Is this always true? What about the bounded-storage model?
- Privacy amplification is secure even against adversaries with quantum knowledge.

This has applications for security proofs of quantum cryptographic schemes.